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(Cin-Ty A. Lee & Don L. Anderson, P. 1141-1156)



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## Editor's note: how and where does continental crust form?

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“The significance of the continental crust on which we live is self-evident, yet our knowledge remains limited on its origin, its way and rate of growth, and how it has acquired the ‘andesitic’ composition from mantle derived magmas” [1]. In this invited contribution, Cin-Ty Lee offers his perspectives. Cin-Ty Lee has been a professor of geochemistry at Rice University in Houston, Texas (USA), since 2002. He received his bachelor’s degree in geology from the University of California, Berkeley in 1996, and did a senior thesis with George Brimhall, an economic geologist. He received his PhD in geochemistry from Harvard University in 2001, working with Roberta Rudnick on the origin and evolution of continental lithospheric mantle. He spent the following year at the California Institute of Technology with Gerald Wasserburg, working on platinum group element geochemistry of marine sediments. He has been a visiting scientist at the Academia Sinica in Taiwan, the Atmosphere Ocean Research Institute in the University of Tokyo, and most recently, he was a Miller visiting professor at UC Berkeley. He is currently secretary of the Volcanology, Geochemistry, and Petrology section of the American Geophysical Union and an editor of Geochemistry, Geophysics, and Geosystems (G-cubed). He has published over 100 peer-reviewed papers, including two first authored publications in *Science* and three in *Nature*. He is the recipient of the American Geophysical Union’s Kuno Award, the Geochemical Society’s Clarke Medal, the Geological Society of America’s Donath medal, and the Packard Fellowship. He is a fellow of the Geological Society of America and the Mineralogical Society

of America. In 2013, he toured China giving the Shen-Su Sun lecture series. He has made important contributions to understanding the origin and evolution of continental crust, continental lithospheric mantle, mantle differentiation, redox processes, and ore formation. Currently, he is investigating the interplays between the deep Earth, oceans, and atmospheres, with a particular focus on understanding long-term climate change and the whole Earth carbon and oxygen cycles. A large component of this work is understanding contact and regional metamorphism, the origin and transport of magmas, and the feedbacks between erosion and magmatism ([www.arc2climate.org](http://www.arc2climate.org)). He has also embarked on a study of the thermodynamics of earthquakes. Lee brings to the table a diversity of toolkits, which includes mass spectrometry, petrography, structural geology, mineralogy, modeling of geodynamic and geochemical processes, and calculation of the physical properties of rocks. In addition to geology ([www.cintylee.org](http://www.cintylee.org)), he has a deep interest in all things related to nature, including teaching and publishing on birds and plants.

I am delighted that Cin-Ty Lee accepted the invitation to contribute an article to *Science Bulletin* on one of the fundamental yet still controversial topics—formation of the continental crust. Lee accepts the standard model that the continental crust is formed at arcs, but offers an innovative perspective to resolve one of the difficulties with the “arc model”, i.e., the primitive arc magma is basaltic, whereas the bulk continental crust is andesitic. He proposes that the arc basalt differentiates into a mafic garnet pyroxenite cumulate (“arclogite”) and a complementary silicic melt. The “arclogite” due to its high density can sink periodically into the convective mantle, whereas the silicic melt becomes the very material for the continental crust. Although this new idea and the standard arc model differ from my own interpretations [1], Lee’s paper revives the

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important debate. Removal of the arclogite would lead to the silicic arc melts (hence, the resulting continental crust) depleted in heavy rare elements (i.e., the “garnet signature”), but the latter is not observed. The arclogite of cumulate origin would have elevated abundances of heavy rare earth elements, depleted in light rare earth elements and other incompatible elements. Hence, the arclogite cannot be the ideal material for hot spot magmatism as proposed. Nevertheless, I am pleased that the model by Lee provides an impetus for enthusiastic scientists to participate in this debate toward a genuine understanding of the origin

and evolution of continental crust—a quest that has attracted the attention of geologists for centuries and the curiosity of the mankind for millennia.

## Reference

1. Niu Y, Zhao Z, Zhu D-C et al (2013) Continental collision zones are primary sites for net continental crust growth—a testable hypothesis. *Earth-Sci Rev* 127:96–110



# Continental crust formation at arcs, the arclogite “delamination” cycle, and one origin for fertile melting anomalies in the mantle

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**Abstract** The total magmatic output in modern arcs, where continental crust is now being formed, is believed to derive from melting of the mantle wedge and is largely basaltic. Globally averaged continental crust, however, has an andesitic bulk composition and is hence too silicic to have been derived directly from the mantle. It is well known that one way this imbalance can be reconciled is if the parental basalt differentiates into a mafic garnet pyroxenitic residue/cumulate (“arclogite”) and a complementary silicic melt, the former foundering or delaminating into the mantle due to its high densities and the latter remaining as the crust. Using the Sierra Nevada batholith in California as a case study, the composition of mature continental arc crust is shown in part to be the product of a cyclic process beginning with the growth of an arclogite layer followed by delamination of this layer and post-delamination basaltic underplating/recharge into what remains of the continental crust. A model is presented, wherein continuous arc magmatism and production of arclogites in continental arcs are periodically punctuated by a delamination event and an associated magmatic pulse every  $\sim 10$ – $30$  My. The recycling flux of arclogites is estimated to be  $\sim 5\%$ – $20\%$  that of oceanic crust recycling by subduction. Delaminated arclogites have the necessary trace-element compositions to yield time-integrated isotopic compositions similar to those inferred to

exist as reservoirs in the mantle. Because of their low melting temperatures, such pyroxenites may be preferentially melted, possibly forming a component of some hotspot magmas.

**Keywords** Pyroxenite · Eclogite · Delamination · Cumulate · Continental crust

## 1 Introduction

Delamination (used loosely here to describe any foundering or detachment) of lower crust or lithospheric mantle due to compositionally or thermally induced densifications has been suggested to explain a number of geologic observations, such as short-lived uplifts, high heat flow and magmatism, and unusual seismic anomalies [1–12]. Such features as uplift and high heat flow are the predicted consequences of having hot asthenospheric mantle upwell passively to replace the “void” created by removal of the deep lithosphere or lower crust. These features have been observed in a number of areas and are increasingly being taken as indirect evidence for recent delamination [11, 13]. If delamination is a general phenomenon, it should be an important means of recycling lower crust or lithospheric mantle back into the Earth’s interior, and hence should have an important influence on the compositional evolution of continental crust and the introduction of compositional heterogeneities into the mantle [5, 9, 14]. For example, it has been hypothesized that the felsic composition of continental crust may be the result of preferential removal of mafic lower crust by delamination [4, 9, 11, 15–30]. It has also been hypothesized that this delaminated mafic reservoir may partly contribute to the source regions of mid-plate and ridge magmas [26, 31].

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There are, however, two endmember ways of generating mafic lower crust. One way is via partial melting of basaltic crust during hot subduction or continental collisions, resulting in the generation of mafic restites that may eventually founder back into the mantle [27, 32–38]. Another way is by crystal fractionation of a basalt, generating mafic cumulates at depth in subduction zone volcanoes [4, 11, 18–23, 29, 30, 39–42]. The purpose of this paper is not to debate the mechanisms by which continental crust is formed. Instead, we focus solely on arc magmatism and evaluate whether delamination of mafic lower crust in arcs is an important geologic process in the evolution of arc crust and the generation of fertile melting heterogeneities within the mantle.

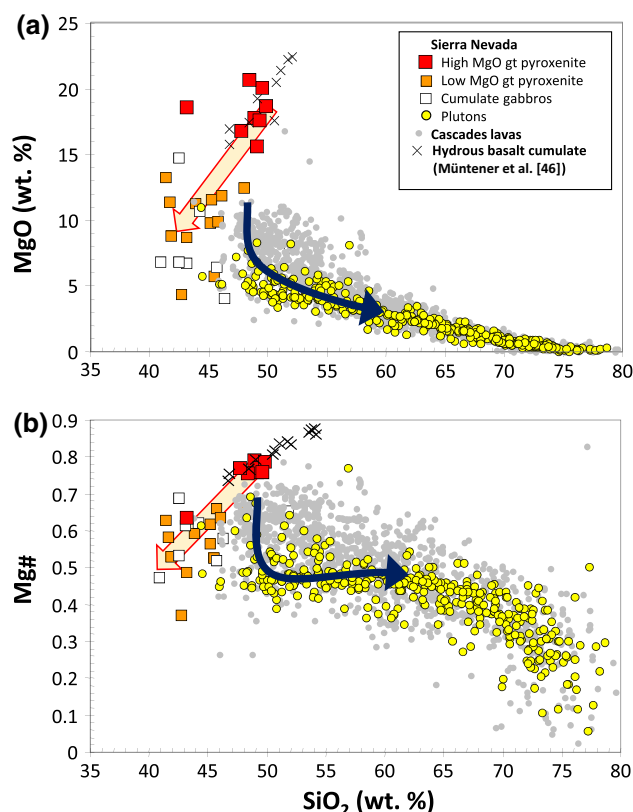
## 2 Delamination in continental arcs

### 2.1 Garnet pyroxenites and their hypothesized delamination

The difficulty in testing the delamination–crust formation hypothesis is that, if it is correct, some of the evidence for delamination is missing, that is, the putative delaminated material has already been disposed of and is presently lying somewhere hidden in the Earth's interior. However, one place where we have been afforded samples of mafic lower crust prior to delamination is in the Sierra Nevada, the eroded remnant of a Mesozoic arc batholith in California [11, 18]. Small-volume basaltic eruptions in the late Miocene (8 Ma) harbor fragments (xenoliths) of a deep and cold mafic garnet pyroxenitic root, while Pleistocene eruptions contain fragments of hot asthenospheric mantle at equivalent depths [11]. This suggests that a thick, cold mafic root existed beneath the central and eastern Sierras up until the late Miocene, but that sometime in the Pliocene was removed wholesale and replaced by asthenospheric mantle [11, 43]. This interpretation is broadly consistent with the observation that, although the highest elevations in the Sierra Nevada are in the east, the underlying crust is thinnest in this region, which implies that the high elevations are supported by hot (hence low density) asthenospheric mantle [10]. Late Pliocene flare-ups of small-volume basaltic magmatism in the central Sierra might even be the manifestation of post-delamination magmatism [44, 45]. Finally, hypothetical delamination in the central and eastern Sierra during the Pliocene may have even propagated westward as evidenced by high seismic velocity anomalies at depth beneath the western edge of the Sierra Nevada; these have been interpreted to represent ongoing convective downwellings of lower crust or lithospheric mantle [13].

### 2.2 The relationship between garnet pyroxenites and Sierran mafic to intermediate plutons

Regardless of whether delamination occurred, it is natural to speculate whether the Sierran garnet pyroxenites are complementary to the Sierran batholith. Ducea and Saleeby [29, 30] showed that the Sierran garnet pyroxenites are isotopically similar to the Sierran granitoids and that their major element compositions are roughly complementary. Their observations hint at a petrogenetic relationship with the Sierran granitoids, but the exact nature of the link was not clear. The link can be clarified by considering here a more comprehensive database of Sierran garnet pyroxenites and Sierran granitoids [18]. The Sierran garnet pyroxenites can be divided into two groups, a high MgO (MgO >13 wt%) and a low MgO group (<13 wt%) as shown in Fig. 1a. The former is characterized by high pyroxene to garnet ratios, high Mg# (molar Mg/(Mg + Fe)), high Ni and Cr contents, and low Al<sub>2</sub>O<sub>3</sub>



**Fig. 1** MgO (a), Mg# (b) versus SiO<sub>2</sub> in Sierran garnet pyroxenites, Sierran and Peninsular Ranges Batholith plutons, and Cascades volcanics, simplified from Lee et al. [18]. Mg# represents molar Mg/(Mg + Fe<sub>T</sub>), where Fe<sub>T</sub> represents total iron. Red outlined arrow represents crystal line of descent, and dark blue arrow represents liquid line of descent. Differentiation at relatively constant Mg# requires precipitation of Fe-bearing phases, such as magnetite, and/or basaltic recharge. X symbols refer to experimentally determined cumulates from a hydrous basalt from Müntener et al. [46]

contents, while the latter is characterized by low pyroxene to garnet ratios, low  $\text{SiO}_2$ , lower Mg#, low Ni and Cr contents, higher  $\text{Al}_2\text{O}_3$  contents and higher heavy rare earth element (HREE) abundances [18] (Figs. 1b, 2a).

An equally important observation is that the Sierran mafic to intermediate magmas have high  $\text{Al}_2\text{O}_3$  (not shown) and low MgO and Mg# (Fig. 1) for a given  $\text{SiO}_2$  content compared to basaltic magmas from mid-ocean ridges or young arcs, such as the Cascades. For example, at  $\text{SiO}_2$  contents of 50 wt% (typical for basalt), Sierran magma MgO contents are anomalously low, requiring an initial decrease in MgO at constant  $\text{SiO}_2$  (Fig. 1a). While the basaltic differentiation trend seen in the Cascades data set can be easily explained by olivine-dominated crystallization, the low MgO and high  $\text{Al}_2\text{O}_3$  of the Sierran mafic plutons cannot. Instead, the mafic to intermediate leg of the Sierran magmatic differentiation series requires the initial removal of pyroxene-rich residues/cumulates (to drive

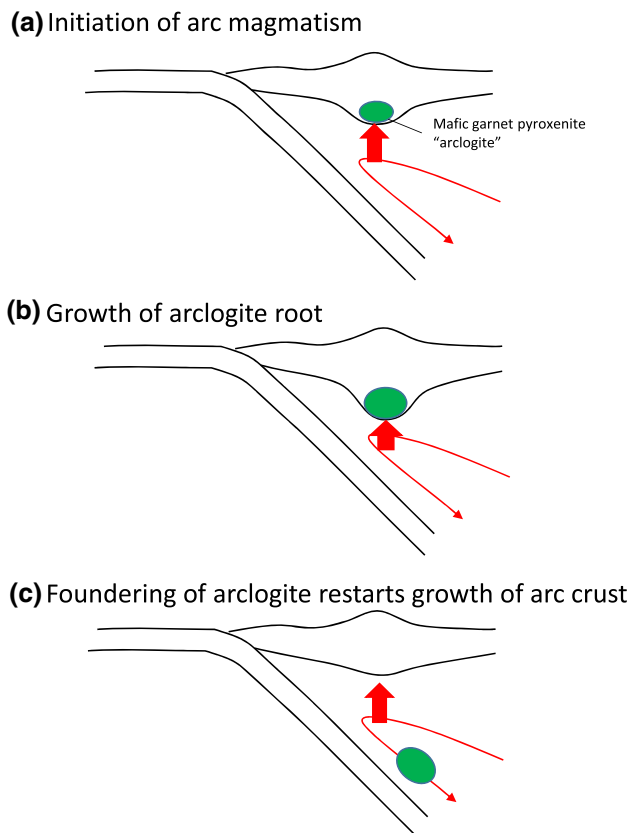
MgO down at near constant  $\text{SiO}_2$ ) followed by removal of  $\text{SiO}_2$ -poor or pyroxene-poor residues/cumulates (to drive low MgO magmas toward higher  $\text{SiO}_2$ ). It can be seen that these hypothesized residue/cumulate compositions are matched by the Sierran high and low MgO garnet pyroxenite compositions. These observations corroborate the petrogenetic link between Sierran garnet pyroxenites and the Sierran mafic to intermediate plutons. To distinguish these garnet pyroxenites from true eclogites, which have omphacitic pyroxene, we refer to them as arclogites, a term first suggested by Anderson [47].

### 2.3 A cumulate origin for Sierran garnet pyroxenites

The origin of the Sierran garnet pyroxenites can be clarified further. The high MgO pyroxenites are not likely to be melt residues because their MgO contents are much too high and would require that they be the residues of >90 % melt extraction [48, 49]. In addition, should the high MgO pyroxenites be the products of melting preexisting basalt, a continuous spectrum of residue compositions would be expected, yet no continuous compositional spectrum is seen in the pyroxenites. The low MgO pyroxenites also cannot represent residues of re-melting primitive basaltic compositions: Melting of eclogitized basalt results in residues with higher MgO and lower  $\text{SiO}_2$  [48, 49], but the low MgO pyroxenites have MgO contents lower than any hypothetical primitive mantle-derived basalt. Thus, the low MgO pyroxenites are likely cumulates of evolved basalt that has already crystallized high MgO-type pyroxenites. This genetic relationship between the two pyroxenite groups is also supported by the fact that Mg#, Ni and Cr contents decrease going from high to low MgO pyroxenites [18]. Interestingly, the low MgO pyroxenites have major element compositions very similar to those of Sierran gabbros, which show unequivocal cumulate textures [50]. The Sierran garnet pyroxenites likely crystallized at high pressures (and possibly wet conditions [46]) in order to preferentially crystallize pyroxene instead of olivine (Fig. 1). If so, the Sierran garnet pyroxenites may represent a semi-continuous series of cumulates that is complementary to the fractionation trend seen in mafic to intermediate (50 wt%–60 wt%  $\text{SiO}_2$ ) Sierran magmas (more silicic plutons, such as granites, are unrelated to these pyroxenites). These interpretations are consistent with those based on modeling of Kohistan pyroxenites and felsic rocks [20, 22].

### 2.4 The need for magmatic recharge

One feature that cannot be explained by fractional crystallization alone is the observation that once the MgO content and Mg# of the Sierran parental magmas have dropped as a consequence of crystallizing high MgO



**Fig. 2** (Color online) Cartoon showing how arc crust magmatically thickens. In **a**, subduction drives corner flow in the mantle wedge (red line), leading to decompression melting (large red arrow). These melts rise and intrude or underplate the over-riding plate, causing the crust and associated deep crustal mafic cumulate pyroxenites (green) to thicken (**b**). Dense pyroxenite layer eventually reaches a critical thickness after which it founders into the mantle wedge (**c**). Continued melting in the mantle wedge, provided subduction continues, re-initiates the cycle of arc crust growth

pyroxenites, the Mg# remains relatively constant with increasing SiO<sub>2</sub> despite our suggestion that this differentiation trajectory is controlled by crystallization of low MgO pyroxenites. Low MgO pyroxenites have higher Mg#s than most of the Sierran plutonic compositions, and thus, fractional crystallization should decrease magmatic Mg#. The constancy of Mg# in the mafic to intermediate spectrum of Sierran plutons indicates that Mg# is buffered (Fig. 1b). One possibility is that this is due to the fractionation of more Fe-rich minerals, such as garnet or cumulate assemblages containing magnetite, giving rise to the well-known calc-alkaline differentiation trend seen in subduction zone magmas [51]. Another explanation, not necessarily exclusive, is that fractional crystallization of garnet pyroxenites is accompanied or followed by basaltic recharge of the residual magma in the form of simultaneous crystallization and recharge of a magma chamber and/or by incremental magmatic underplating/mixing with the lower crust after each magmatic differentiation event [18, 52].

### 3 The arclogite delamination cycle in continental arcs

Based on the above discussion, garnet pyroxenite accumulation may be a fundamental process in the formation of mature continental arcs, that is, arcs built on preexisting continental lithosphere. In such cases, mantle wedge-derived magmas must pass through a thicker lithosphere and therefore are likely to begin crystallization at greater depths than what might be seen beneath island arcs or young, incipient continental arcs (e.g., the Cascades). The higher pressures of crystallization will favor the precipitation of pyroxenes over olivine.

We propose the following model for the formation and evolution of mature continental arcs (Fig. 2). Our model is constructed to satisfy (1) the dominantly cumulate origin of Sierran garnet pyroxenites, (2) the need for basaltic recharge/underplating, and (3) the possibility that the garnet pyroxenites have delaminated. Thus, as an arc matures, a garnet pyroxenite layer begins to build up at its base due to the progressively thicker lithosphere through which arc magmas must pass (Fig. 2b). Given a continuous background magmatic flux imparted by subduction-induced decompression melting in the mantle wedge [53], the garnet pyroxenite layer thickens gradually with time [4, 18, 54]. Due to the high densities of the garnet pyroxenite cumulates compared to typical peridotitic mantle (Fig. 3; [55]), the garnet pyroxenite layer will eventually founder/delaminate once a critical thickness is reached (Fig. 2c). Hot, asthenospheric mantle should then rise passively to fill the void generated by delamination and, in so doing, generate a pulse of magmatism superimposed on the background flux of melting associated with decompression in

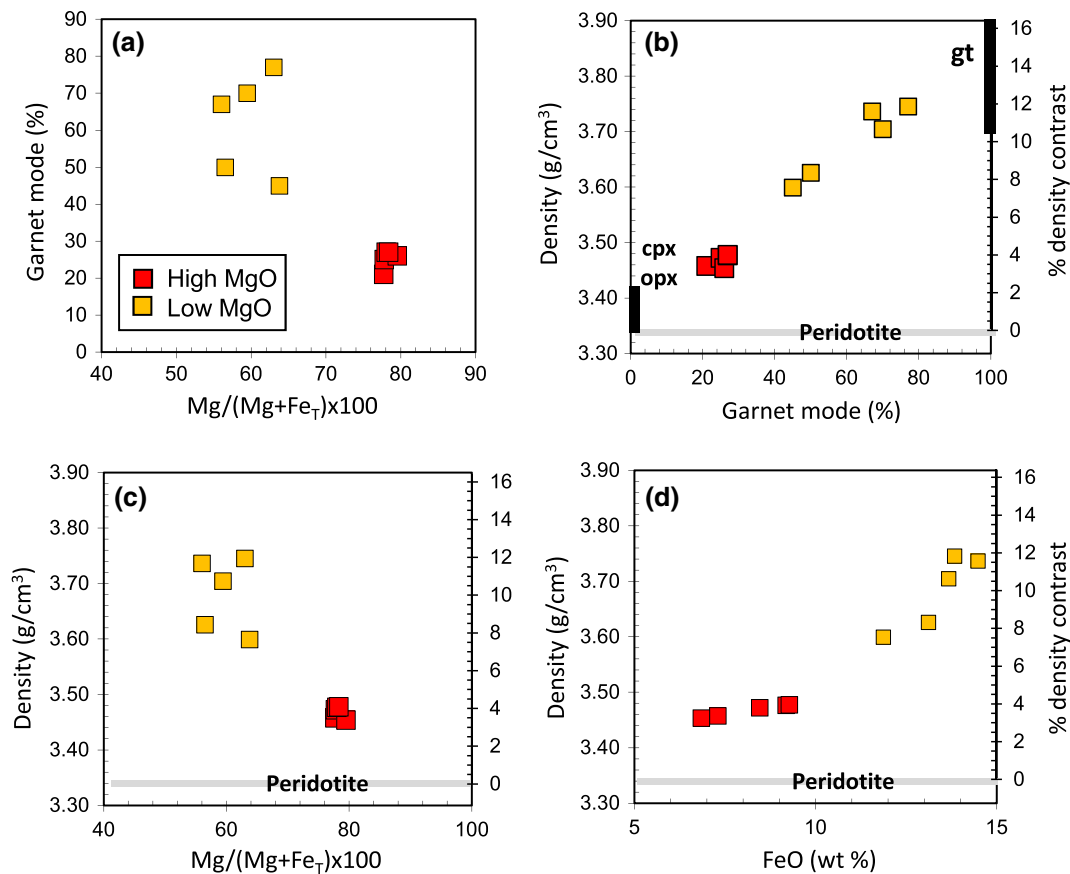
the mantle wedge. This renewed melting results in basaltic underplating of the more felsic crust. Some of the basalt will mix with and heat up the base of the overlying continental crust, causing it to melt and generate more evolved magmas, while some basalt will simply crystallize to form more garnet pyroxenite cumulates. If subduction-related magmatism continues, the cycle of growing a cumulate layer followed by delamination will reinitiate. In this way, Sierran magmas can evolve to intermediate compositions without a significant change in Mg# (apart from the initial decrease associated with crystallization of olivine-bearing lithologies and high MgO pyroxenites).

In Fig. 4, we plot the compositions of Sierran pyroxenites, felsic plutons and estimates of the composition of the continental crust relative to mid-ocean ridge basalt (MORB). It can be seen in Fig. 4a that Sierran felsic plutons and the global bulk continental crust, including the present-day global lower continental crust, are depleted in CaO, FeO and MgO relative to basalt. Cumulates with high CaO, FeO and MgO are needed for mass balance, but the present-day lower continental crust clearly does not have the appropriate composition. Arc pyroxenite cumulates, as can be seen from Fig. 4b, have the desired complementary composition. In summary, these qualitative mass balance considerations indicate that the formation of felsic magmatic arcs must in general be accompanied by a significant return of mafic cumulates back into the mantle.

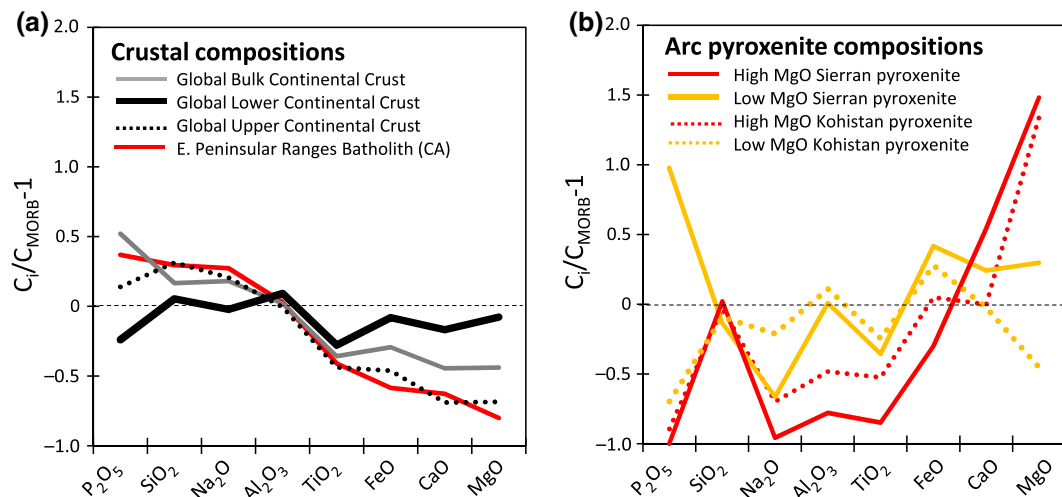
## 4 Simple models and predictions

### 4.1 Estimating cumulate growth rate and delamination flux in arcs

Assuming our conceptual model in Sect. 3 is correct, we are now left with quantifying two parameters: (1) the magnitude of the delamination flux of mafic garnet pyroxenites, and (2) the average periodicity of delamination. The delamination flux can be estimated by first considering a mass balance between the two pyroxenite groups and average Sierran pluton compositions with respect to primitive arc basalt. Based on inversion of major element oxides, we have shown that the Sierran garnet pyroxenites collectively amount to 50 %–70 % by mass of the parental basaltic magma, as shown in Fig. 5 [4]. To convert this number to a cumulate growth rate requires that we know the Sierran magmatic flux during the Mesozoic. Because we do not know this, we assume for simplicity that arc magmatic fluxes are relatively uniform globally, regardless of whether mature continental arcs or island arcs are considered. We can then take 3–9 km<sup>3</sup>/year as the average total production rate of magmas in arcs [61, 62], ~51,000 km as the total length of subduction zones [57],

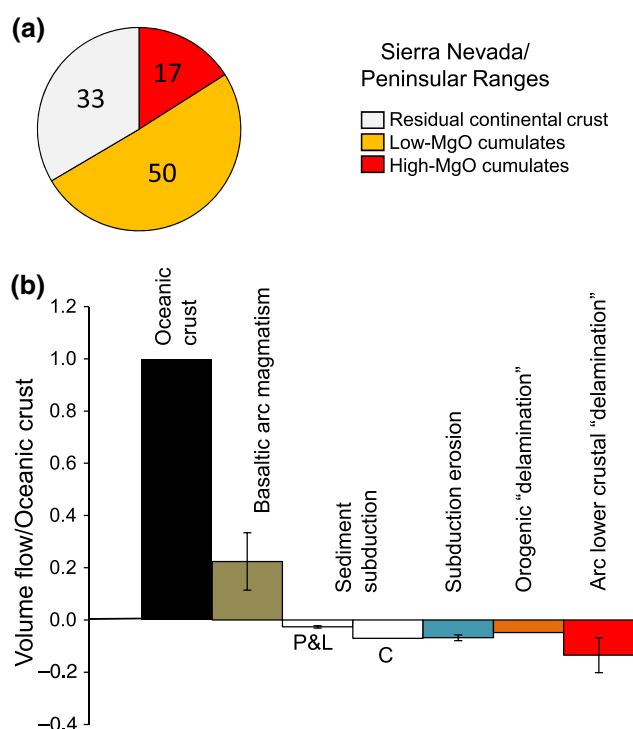


**Fig. 3** (Color online) **a** Garnet mode versus Mg# (molar  $\text{Mg}/(\text{Mg} + \text{Fe}_T) \times 100$  where  $\text{Fe}_T$  represents total Fe) in Sierran arc garnet pyroxenites, **b** density in  $\text{g}/\text{cm}^3$  versus garnet mode, **c** density versus Mg#, **d** density versus bulk FeO (total Fe). Density contrast relative to peridotite is shown on the right-hand y-axis (peridotite density taken from Lee [55]). All densities calculated at standard state and pressure (STP) conditions. Figures adapted or using data from Lee et al. [18]



**Fig. 4** (Color online) Crustal type **(a)** and arc pyroxenite **(b)** major and minor element compositions using data from Lee [4] and Lee et al. [19] for Sierra Nevada and Jagoutz et al. [20, 22, 39] for Kohistan arc. Average composition of mid-ocean ridge basalt (MORB) is from Arevalo and McDonough [56]. All compositions have been normalized to MORB and represented as relative deviations. Thus, MORB plots at zero, concentrations higher than MORB plot positive and concentrations lower than MORB plot negative. In **(a)** it can be seen that global average lower continental crust cannot balance the composition of the upper continental crust. In **(b)** it can be seen that arc pyroxenites are mafic enough to balance both the composition of the upper continental crust and the bulk continental crust relative to a basaltic parent





**Fig. 5** (Color online) **a** Mass proportions of mafic pyroxenites (low and high MgO) and felsic residual magmas resulting from differentiation of a primary mantle-derived arc basalt. Values are taken from Lee [4] based on inversion of major element compositions, **b** volume flow rates (km<sup>3</sup>/year) of various magmatic or sedimentary fluxes, normalized to modern oceanic crust production rates (22.9 km<sup>3</sup>/year assuming 7-km-thick crust and areal oceanic crust production rate from Bird [57]). Volume flow rates for arc magmatism, sediment subduction [58, 59], subduction erosion [60], orogenic delamination [58], and arc cumulate delamination [4] are also shown

and ~100 km as the width of magmatic arcs, to arrive at 0.3–1.2 km/My for the average rate of crustal thickening in an arc by magmatic inflation. This corresponds to a cumulate thickening rate ~0.2–0.8 km/My if the mafic cumulates are assumed to be 65 % of the parental magma, as discussed above [4]. Given the higher densities of the pyroxenites compared to felsic plutons and erupted basalts, these numbers might need to be reduced by ~15 %. However, given the uncertainties in all the flux estimates, we do not make this correction in our results. Our estimated arc and delamination fluxes are shown in Fig. 5b normalized to average oceanic crust production rates today (22.6 km<sup>3</sup>/My; [57]). Volume flow rates for arc magmatism, sediment subduction [58, 59], subduction erosion [60], orogenic delamination [58], and arc cumulate delamination [4] are also shown for comparison. It can be seen that delamination rates of arc lower crust, while much smaller than the subduction rate of oceanic crust, is comparable if not larger than the amount of sediment being subducted.

## 4.2 Delamination versus viscous foundering of the pyroxenite layer

We now attempt to predict the periodicity of episodic delamination events by considering how long it takes for the cumulate layer to reach critical negative buoyancy beyond which foundering takes place. The tendency to founder is controlled by the competing effects of negative buoyancy forces (associated with compositional and thermal densification and the thickness of the cumulate layer) and viscous resisting forces. Thermobarometry of Sierran garnet pyroxenites indicates that prior to their hypothesized removal in the Pliocene, they last equilibrated at temperatures below ~800 °C at ~50 km depth [11, 43, 63–66]. This suggests that, following the formation of the cumulate layer, considerable cooling took place, most likely due to impingement upon the cold subducting Farallon plate as the arc root thickened [66].

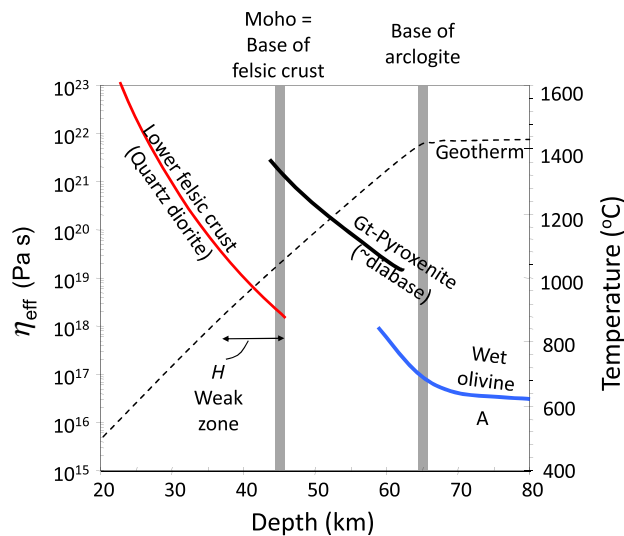
Using a diabase rheology as an analog for the rheology of garnet pyroxenite [67], we find that the effective viscosity at these temperatures is ~10<sup>22</sup> Pa s (Table 1; Fig. 6). This is probably a minimum estimate because garnet pyroxenite is probably stronger than diabase. These high viscosities imply that a garnet pyroxenite cumulate layer will be effectively unable to founder viscously even if it is denser than the underlying mantle (Fig. 6). Only at higher temperatures (1,200 °C) can viscous foundering of the pyroxenite layer occur and the problem can be treated as a Rayleigh–Taylor instability [5]. In the case of the Sierra Nevada, the pyroxenite layer is cold and thus too strong to founder viscously, yet all evidence from seismology and temporal changes in xenolith demographics and basalt thermobarometry indicate that the pyroxenite root was removed [11, 44, 45, 68–70].

A dense but effectively rigid cumulate layer can nevertheless be removed if the layer is separated from the overlying crust by a narrow low-viscosity zone, e.g., a weak lower crustal layer. Taking quartz diorite rheology [65] as an analog for the lower part of the more felsic part of the continental crust lying above the garnet pyroxenite layer, a lower crustal weak layer having effective viscosities of ~10<sup>19</sup> Pa s and a width of roughly 1–5 km is obtained (Fig. 6). This weak zone allows for wholesale removal of the cumulate layer as required by the

**Table 1** Flow parameters for some lithospheric materials

|                              | Granite | Wet granite | Diabase           | Olivine            |
|------------------------------|---------|-------------|-------------------|--------------------|
| $A_0$ (GPa <sup>-n</sup> /s) | 5       | 100         | $3.2 \times 10^6$ | $4 \times 10^{15}$ |
| $n$                          | 3.2     | 1.9         | 3.4               | 3                  |
| $E_A$ (kJ/mol)               | 123     | 137         | 260               | 540                |

Flow law parameters are taken from [67]

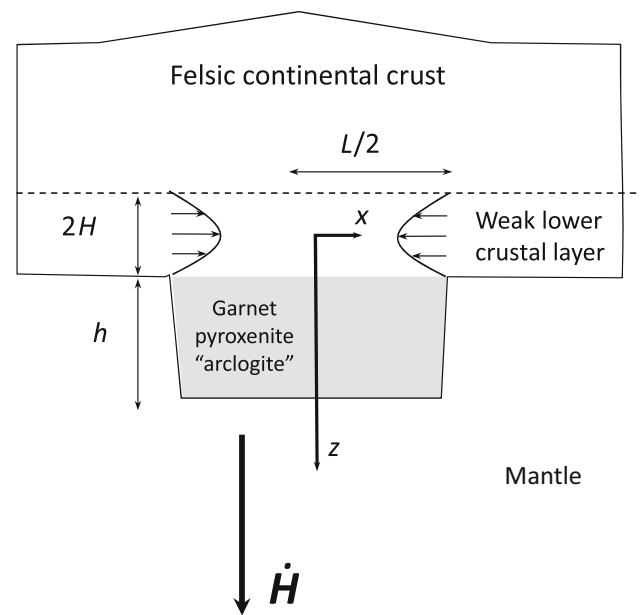


**Fig. 6** This figure shows effective viscosities on the y-axis versus depth for a hypothetical layered lithosphere (felsic crust, underlain by garnet pyroxenite, e.g., arclogite layer and mantle peridotite). Effective viscosities are calculated using Eq. 3 and the flow laws are shown in Table 1 along a hypothetical geotherm for an arc (dashed line with the right-hand y-axis for reference). A background stress state of 0.3 MPa was assumed. Note the presence of a weak lower crustal layer just above the interface between the felsic and mafic (pyroxenite) crust

interpretation of xenolith data in the Sierran case study (see Sect. 2.1). This rheological stratification is idealized in Fig. 7 for the purposes of modeling in the next section. Whether or not the dense cumulate layer founders on geologically reasonable timescales will depend on how large the viscous resisting forces are in the weak crustal layer. As the pyroxenite layer sinks, it will suck material into the thin gap, generating a low-pressure region that resists deformation and slows down delamination. The viscous resisting forces in the weak gap will be greatest when the gap is thin. Detachment results in the widening of the gap until viscous resisting forces within the gap become negligible compared with viscous resistance from the underlying asthenospheric mantle.

#### 4.3 Analytical model for delamination

There are a number of sophisticated numerical models describing the detachment of dense mafic lower crust [5, 71–73]. The approach here is to develop simple analytical models in the spirit of Bird [1]. Our main purpose for developing an analytical solution is to make it easier to isolate variables. In particular, we are interested in the case in which the garnet pyroxenite layer can grow with time via magmatic underplating. The basic conceptual approach along with models for simple Newtonian fluids was developed in a pedagogical manner in Lee [4]. Here, we



**Fig. 7** Highly idealized cartoon of the model for detachment of the pyroxenite (arclogite) layer. The pyroxenite layer is assumed to be separated rheologically from the upper crust by a weak lower crustal layer/gap (Fig. 6). Because the pyroxenite layer is dense, it will tend to sink, resulting in the inward flow of lower crustal material. This lower crustal flow is approximated by laminar flow in the foregoing model

consider the more general case of non-Newtonian rheologies. The model geometry is shown in Fig. 7. Variables and their units are shown in Table 2.

The general flow law for a non-Newtonian rheology is given by

$$\dot{\epsilon}_{ij} = A \sigma_E^{(n-1)} \sigma_{ij}, \quad (1)$$

where  $\sigma_E$  is the second invariant of the deviatoric stress tensor,  $A$  is related to viscosity according to  $A = A_0 e^{-E_A/RT}$ ,  $R$  is the gas constant,  $T$  is temperature (Kelvin),  $n$  is the power-law exponent,  $A_0$  is the preexponential factor with units of  $G \text{ Pa}^{-n}/s$ , and  $E_A$  is the activation energy with units of  $\text{kJ/mol}$  [74, 75]. We ignore the pressure dependence of creep because the variation in lithostatic pressures in our problem is small. Using the following relation

$$\dot{\epsilon}_E = A \sigma_E^n, \quad (2)$$

the effective viscosity can be determined

$$\eta = A^{-1/n} \dot{\epsilon}_E^{(1-n)/n}. \quad (3)$$

If the width of the weak or low-viscosity layer  $H$  is small compared to the horizontal extent of the cumulate layer  $L$ , the motion of the fluid in the thin gap can be approximated as being parallel to the upper and lower boundaries of the gap (e.g., a lubrication approximation is taken). If flow is laminar only in the  $x$  direction, the following holds

**Table 2** List of parameters and units

|                       | Units       | Description                                     |
|-----------------------|-------------|---|
| $\dot{\epsilon}_{zx}$ | $s^{-1}$    | Strain rate in direction of $x$ , normal to $z$ |
| $\sigma_{zx}$         | Pa          | Shear stress                                    |
| $\dot{\epsilon}_E$    | $s^{-1}$    | Strain rate invariant                           |
| $\sigma_E$            | Pa          | Stress invariant                                |
| $A_0$                 | $Pa^{-n}/s$ | Preexponential factor in flow law               |
| $\eta$                | Pa s        | Effective viscosity                             |
| $n$                   |             | Power in flow law                               |
| $P$                   | Pa          | Pressure  |
| $Q$                   | $m^2/s$     | Volume flux per unit width                      |
| $E_A$                 | kJ/mol      | Activation energy                               |
| $t$                   | s           | Time  |
| $L$                   | m           | Length of pyroxenite layer                      |
| $H$                   | m           | Half thickness of low-viscosity layer           |
| $H_0$                 | m           | Initial half thickness of low-viscosity layer   |
| $\Delta\rho$          | $kg/m^3$    | Density contrast                                |
| $T$                   | $^{\circ}C$ | Temperature                                     |
| $h$                   | km          | Thickness of pyroxenite layer                   |
| $\dot{h}$             | km/year     | Growth rate of pyroxenite layer                 |
| $v$                   | m/s         | Velocity of fluid in slab gap                   |
| $\dot{H}$             | m/s         | Sinking velocity                                |
| $F$                   | $kg/s^2$    | Force per unit width                            |

$$\dot{\epsilon}_E = \left[ \frac{1}{2} (\dot{\epsilon}_{ij} \dot{\epsilon}_{ij}) \right]^{1/2} = \dot{\epsilon}_{zx}, \quad (4)$$

where  $\dot{\epsilon}_{zx}$  is the shear strain rate in the horizontal direction  $x$  within the weak lower crustal layer. Thus, Eq. 2 becomes

$$\sigma_{zx} = -\frac{1}{A^{1/n}} \dot{\epsilon}_{zx}^{1/n}, \quad (5)$$

which describes shear stresses in the weak lower crustal layer. We will assume for now that viscous resistance in this lower crustal weak layer is initially much higher than that in the underlying asthenospheric mantle, such that we can approximate the onset of sinking of the pyroxenite layer by considering viscous resistance in the weak layer only. Then, assuming that the lateral pressure gradient in the weak crustal layer is constant, the equations of motion can be integrated to yield

$$\sigma_{zx} = -z \frac{\partial P(x)}{\partial x}. \quad (6)$$

Assuming that driving forces and viscous resistance forces balance (a zero inertia system is assumed), these two equations can be equated and integrated with respect to  $z$  to yield an expression for the lateral fluid velocity as a function of gap width  $H$

$$v_x = \frac{A}{n+1} \left( \frac{\partial P}{\partial x} \right)^n (z^{n+1} - H^{n+1}). \quad (7)$$

Conservation of mass requires that the mass flow of material into the thin gap at every increment in time is balanced by an increase in  $H$  in response to the sinking cumulate layer, assuming that the pyroxenite slab is rigid [73]. The volume flux into the gap is

$$Q = \int_{-H}^H v_x dz = 2 \int_0^H v_x dz, \quad (8)$$

which corresponds to

$$Q_{\text{total}} = -2A \left( \frac{\partial P}{\partial x} \right)^n \left[ \frac{H^{n+2}}{n+2} \right]. \quad (9)$$

This volume flow must equal

$$Q = -(2\dot{H})(L/2) = -\dot{H}L, \quad (10)$$

where  $\dot{H}$  is the sinking velocity of the cumulate layer. Equations 9 and 10 and isolating for  $\frac{\partial P}{\partial x}$  yield

$$\frac{\partial P}{\partial x} = \left[ \frac{\dot{H}L(n+2)}{2H^{n+2}A} \right]^{1/n}. \quad (11)$$

Integrating with respect to  $x$  and applying the boundary condition,  $P = P_{L/2}$  when  $x = L/2$ , yields

$$P_{L/2} - P(x) = \left[ \frac{\dot{H}L(n+2)}{2H^{n+2}A} \right]^{1/n} \left( \frac{L}{2} - x \right). \quad (12)$$

The suction force per unit width (in third dimension),  $F_P$ , is

$$F_P = \int_0^{L/2} (P_{L/2} - P(x)) dx. \quad (13)$$

Integrating (12) yields the total viscous resisting force acting within the weak lower crustal layer

$$F_P = \left[ \frac{\dot{H}^{1/n}}{H^{(n+2)/n}} \right] \left[ \frac{L(n+2)}{2A} \right]^{1/n} \frac{L^2}{8}. \quad (14)$$

Because our system can be assumed to be a zero inertia system (Reynold's number is zero), the viscous resisting force is balanced exactly by the buoyancy force

$$F_g = (\Delta\rho)(L/2)hg, \quad (15)$$

where  $\Delta\rho$  is the density contrast between the pyroxenite layer and the underlying asthenospheric mantle,  $h$  is the thickness of the cumulate layer, and  $g$  is gravity. Equations 14 and 15 give the sinking rate as a function of gap width

$$\dot{H} = \frac{F_g^n}{K^n} H^{n+2}, \quad (16a)$$

$$K = \left[ \frac{L(n+2)}{2A} \right]^{1/n} \frac{L^2}{8}. \quad (16b)$$

Equation 16a shows that the sinking velocity increases with gap width  $H$  to the  $n+2$  power and therefore increases with time. We can integrate Eq. 16a to get  $H$  as a function of time

$$H(t) = \left[ \frac{1}{H_0^{n+1}} - \frac{(n+1)F_g^n}{K^n} t \right]^{-1/(n+1)}. \quad (17)$$

For a cumulate layer of thickness  $h$ , the critical time for foundering to take place occurs when  $H/H_0 \gg 1$  and is expressed as

$$t_{\text{crit}} = \frac{K^n}{F_g^n(n+1)H_0^{n+1}}. \quad (18)$$

We can also define the time at which the gap width increases by twice the initial gap width

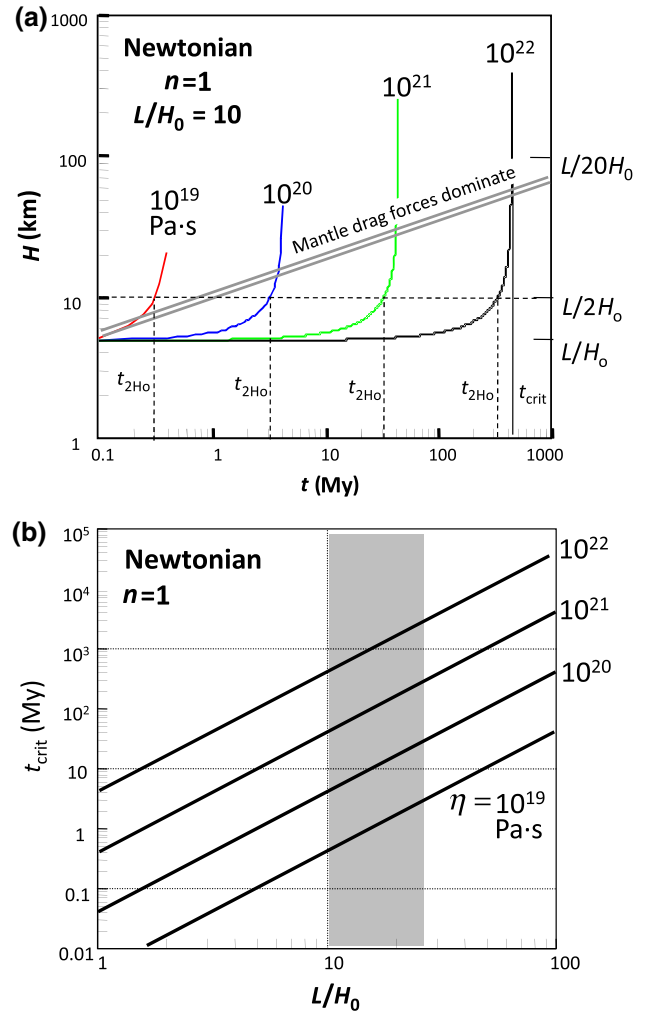
$$t_{2H_0} = \frac{K^n}{(n+1)F_g^n} \left( \frac{2^{n+1}-1}{(2H_0)^{n+1}} \right) = \left( 1 - \frac{1}{2^{n+1}} \right) t_{\text{crit}}. \quad (19)$$

At times greater than  $t_{2H_0}$ , it can be seen that the sinking velocity begins to “accelerate”, in which case, we can assume that delamination is well on its way.

In all of the above equations, if  $n$  is assumed to be 1, the foregoing simplifies to the equations derived by Lee [4] for Newtonian rheologies. Specifically, the time for delamination for  $n=1$  is

$$t_{\text{crit}} = \frac{3}{16} \frac{\eta}{\Delta \rho g h} \left( \frac{L}{H_0} \right)^2, \quad (20)$$

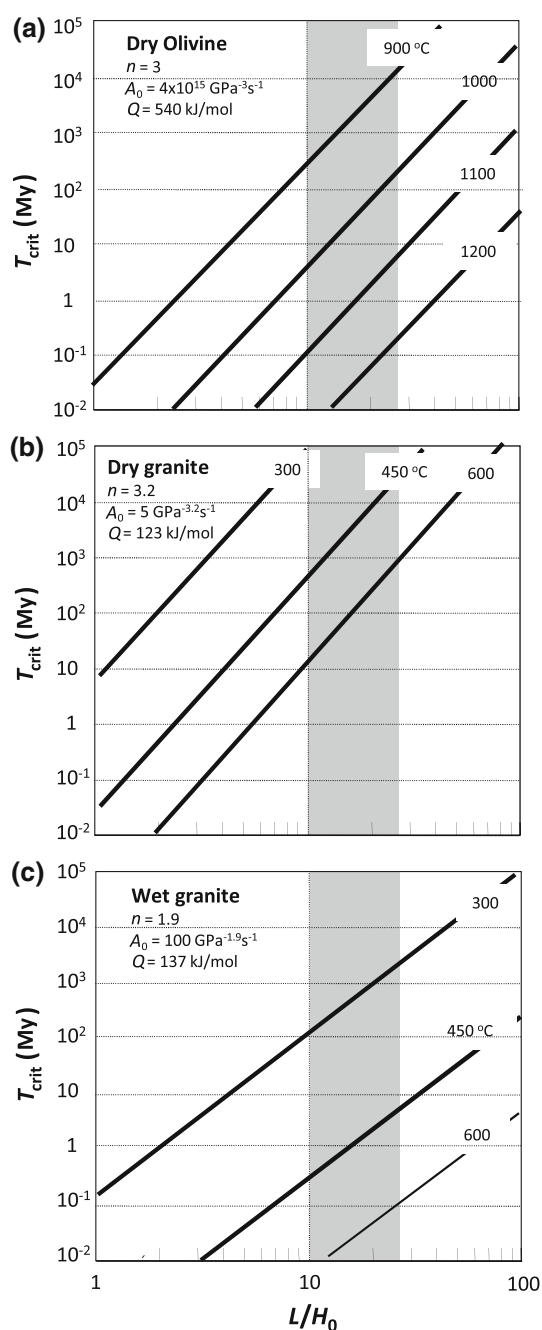
where  $\eta$  is the effective viscosity ( $\eta = 1/A$  for  $n=1$ ). It can be seen that the time for delamination increases with increasing effective viscosity of the lower crustal layer and increasing aspect ratio of the gap  $L/H_0$  and decreasing density contrast and size of the pyroxenite layer (Fig. 8a). Thus, for  $\Delta \rho$  of 500 kg/m<sup>3</sup>,  $h$  of 30 km, and  $\eta$  between  $\sim 10^{19}$  and  $10^{21}$  Pa s, delamination occurs within 100 My for relevant  $L/H_0$ . When  $n > 1$ , as in the case for dislocation creep ( $n \sim 3$ ), the time for delamination will be less than that given for the Newtonian ( $n=1$ ) regime, assuming all other parameters in the flow law are the same. In Fig. 9, we show delamination time  $t_{\text{crit}}$  for non-Newtonian rheologies in which dry olivine, dry granite and wet granite make up the low-viscosity layer, assuming the geotherm shown in Fig. 6. Obviously, the layer is not composed of olivine, so this is shown only for comparison. For a dry granite at 600 °C, delamination occurs before 100 My, and for a wet granite rheology, delamination occurs in less than 1 My. Most likely, the rheology of the weak lower crustal layer will be intermediate between a wet and dry granite, and we conclude that delamination of



**Fig. 8** **a** Delamination,  $t_{\text{crit}}$ , as a function of  $L/H_0$ , showing the dependence of delamination on the aspect ratio of the low-viscosity gap. For the Sierra Nevada,  $L/H_0$  ranges between 10 and 25 (gray vertical bar), **b** one-half of the gap thickness ( $H$ ) as a function of time (My), where  $H_0 = 5$  km,  $L = 50$  km, and density contrast between pyroxenite slab and asthenosphere  $\Delta \rho = 500$  kg/m<sup>3</sup>, pyroxenite layer thickness  $h = 30$  km, and a Newtonian rheology ( $n=1$ ) with viscosities  $\eta$  ( $=1/A$ ) of  $10^{19}$  to  $10^{22}$  Pa s for the low-viscosity layer.  $H$  grows very slowly initially, dropping off exponentially at  $t_{2H_0}$  (vertical dashed lines), when  $H = 2H_0$  (horizontal dashed line). Delamination time is denoted by  $t_{\text{crit}}$  (e.g., when  $H$  approaches infinite). Diagonal double line represents the point at which asthenospheric mantle viscous resistance to sinking  $F_A$  becomes larger than the viscous resistance in the lower crustal weak layer  $F_P$  (e.g., when  $F_A > F_P$ ), assuming an asthenospheric viscosity of  $10^{21}$  Pa s. As can be seen, exponential growth of  $H$  occurs before  $F_A$  becomes important

the pyroxenite layer should occur spontaneously on time-scales of a few to tens of millions of years.

As noted from the outset, we did not account for the viscous resisting force,  $F_A$ , associated with the underlying asthenospheric mantle (this was also ignored by [1]). Accounting for this effect should increase  $t_{\text{crit}}$ , but by how much? We can estimate the relative contribution of  $F_A$  relative to  $F_P$  from the relation  $\dot{\epsilon}_{xz} = \dot{H}/L$ , where the



**Fig. 9** Effect on  $t_{\text{crit}}$  for non-Newtonian rheologies. **a** Dry olivine for the asthenospheric mantle, **b** dry granite for the upper crust, **c** wet granite for the upper crust. Vertical gray bar represents the range of  $L/H_0$  relevant to the Sierra Nevada. Rheologic laws are taken from Table 1 [67].  $\Delta\rho = 500 \text{ kg/m}^3$  and pyroxenite layer thickness  $h = 30 \text{ km}$

deformation imposed on the surrounding asthenosphere by the sinking slab is  $\dot{\epsilon}_{xz}$ , the velocity of the sinking slab is  $\dot{H}$  (by differentiation of Eq. 9), and the length scale over which deformation in the asthenosphere occurs is taken to scale with the length of the slab,  $L$ . For a Newtonian fluid

( $n = 1$ ), the asthenospheric viscous resisting force (per unit width),  $F_A$ , is then expressed as:

$$F_A = \sigma_{xz}(h + L/2) = \eta \dot{\epsilon} = \frac{\dot{H}}{A_A} \frac{h + L/2}{L}, \quad (21)$$

where the stress is taken to act on the surface area exposed to the asthenosphere ( $h + (L/2)$ ) and  $1/A_A$  is used to denote the Newtonian viscosity of the asthenospheric mantle, which is not the same as that in the low-viscosity gap  $A$ .  $F_A$  can be compared to  $F_P$  for the case in which  $n = 1$ :

$$F_P = \frac{\dot{H}}{A} \left( \frac{L}{H} \right)^3 \frac{3}{16}, \quad (22)$$

$$\frac{F_P}{F_A} = \frac{3}{16} \frac{A_A}{A} \frac{L}{h + (L/2)} \frac{L^3}{H^3}. \quad (23)$$

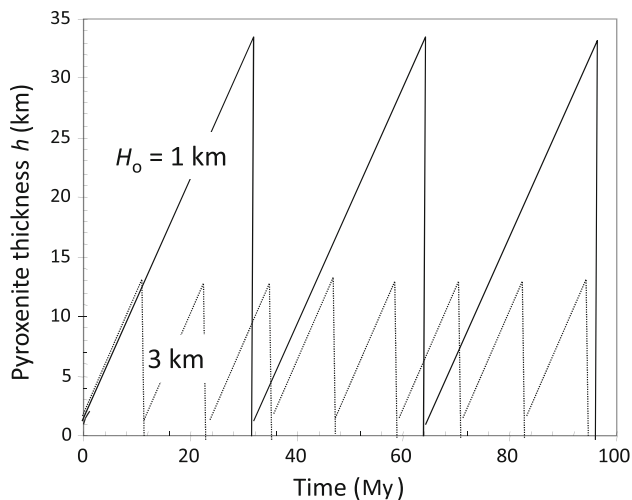
Asthenospheric viscous resistance  $F_A$  becomes important when  $F_P/F_A < 1$ . For a given  $A_A/A$ ,  $F_P/F_A$  decreases with  $1/H^3$ , which is to say that viscous resistance in the gap dominates initially and only after the pyroxenite slab detaches ( $H > 2H_0$ ) does resistance from the asthenospheric mantle become important. In other words, Eqs. 16 and 17 are only valid up until  $F_P/F_A = 1$ . Assuming parameters relevant for the Sierra Nevada ( $L/h = 50 \text{ km}/30 \text{ km}$ ) and the unusual case in which  $A_A = A$ , we find that  $F_P/F_A = 1$  when  $H = 30 \text{ km}$  (Fig. 8b).

We note that the viscosity of the lithospheric mantle may also play a role as in the case of destabilization of cratonic mantle [76–82]. In the case of cratonic lithosphere, the lithospheric mantle could be  $\sim 200 \text{ km}$  thick, whereas the lower crust is less than  $20 \text{ km}$  thick. In the case of arc lithosphere, it seems likely that the lithospheric mantle is thin, given that mature volcanic arcs may have crustal roots extending down to depths in excess of  $60 \text{ km}$  [11, 21, 29, 30, 64–66]. Additionally, in juvenile or island arcs, thermobarometric constraints on the origin of basaltic magmas suggest the presence of hot asthenospheric mantle at depths shallower than  $60 \text{ km}$  [83, 84], leaving little room for lithospheric mantle when typical arc crusts are  $\sim 20\text{--}30 \text{ km}$  thick.

#### 4.4 Application to the general case in which the pyroxenite layer is growing

In the more general case of constant background magmatism, the cumulate layer should be growing by magmatic underplating. If so,  $h$  is a function of  $t$ , and hence,  $F_g$  increases while sinking is taking place. In such a scenario, the above equations are solved incrementally assuming some knowledge of the cumulate growth rate  $\dot{h}$ , which is some fraction of the total juvenile arc magmatic production rate (Fig. 5). Modeling results are shown in Fig. 10 assuming a temperature of  $800^\circ\text{C}$  in the lower crust, a





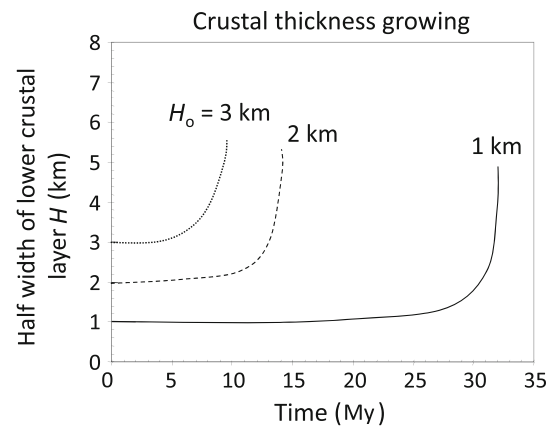
**Fig. 10** Pyroxenite layer thickness  $h$  grows with time at a rate defined by juvenile arc magmatic flux ( $\text{km}^3/\text{km}^2/\text{My}$ ) as represented by the slope of  $h$  versus  $t$ . As the dense pyroxenite layer grows, the negative buoyancy driving force increases as  $h$ , causing the pyroxenite layer to sink and eventually delaminate. We assume wholesale delamination. After delamination, continual magmatic flux permits the pyroxenite layer to grow again, reinitiating the cycle. Calculations are done for different initial thicknesses of lower crustal weak layer (half widths of 1 and 3 km).  $\Delta\rho = 500 \text{ kg/m}^3$  and a power-law rheology of quartz diorite at  $800^\circ\text{C}$  is assumed

quartz diorite rheology for the weak lower crust, a density contrast of  $\sim 500 \text{ kg/m}^3$  based on density estimates for garnet pyroxenites and peridotites [18, 47] (Fig. 3), and a constant cumulate growth rate  $\dot{h}$  (the slope of the line in Fig. 10). It can be seen that for a reasonable range in initial gap widths,  $H_0$ , it takes 10–30 My for foundering to take place (Fig. 11); the cumulate thickness at the time of foundering ranges between 10 and 35 km (Fig. 10), resulting in large heterogeneities recycled into the mantle at any given time. Once the cumulate layer founders, it is replaced by hot asthenosphere, which undergoes decompression partial melting. This results in a magmatic pulse superimposed on a background magmatic flux associated with flux melting in the mantle wedge.

## 5 Implications and predictions

### 5.1 Global implications

In summary, the silicic nature of bulk continental crust can in part be explained by the refinement of basalt at mature continental arcs through a cyclic process of fractional crystallization, delamination, and post-delamination magmatism, all of which are superimposed on a constant baseline magmatic flux associated with flux melting in the mantle wedge. Our conclusions do not rule out other proposed mechanisms for generating felsic continental crust,



**Fig. 11** Detachment rate as a function of time. Same as in Fig. 10 except the effect of initial thickness  $H_0$  of lower crustal weak layer is shown. Crustal thickening rate same as in Fig. 10.  $\Delta\rho = 500 \text{ kg/m}^3$  and a power-law rheology of quartz diorite at  $800^\circ\text{C}$  is assumed

such as partial melting of subducting oceanic crust or over-thickened basaltic crust [16, 17, 32, 34, 35, 38, 85], remelting of sediments [86, 87] or relamination of sediments [88]. Additional work is needed to assess which of these proposed mechanisms dominates the formation of continental crust throughout Earth's history. The immediate prediction of our hypothesis is that continental arc magmatism should be punctuated by magmatic pulses associated with delamination of mafic cumulates, e.g., arclogites. This in turn predicts that continental arcs should be characterized by multiple delamination events. There is some evidence that magmatism in continental arcs, such as the Sierra Nevada, may be episodic [89]. There is also some tentative evidence for multiple delamination events in the Sierra Nevada [43, 90] and the Andes in South America [91]. It is even likely that delamination events in continental arcs are closely coupled with foreland fold-and-thrust belts because removal of arc lower crust and lithospheric mantle provides room for continental basement to be underthrust beneath the arc [92]. Indeed, continental upper plate lithologies have been found at 45 km depth in the Sierran arc, indicating that mafic magmatism and tectonic underthrusting were coeval [93].

In any case, an important implication of our hypothesis is that continental crust formation, at least in continental arcs, is coupled with the recycling of mafic garnet pyroxenites back into the mantle. The question then is how significant is this recycling flux compared to subduction recycling of oceanic crust? Assuming that garnet pyroxenite formation occurs in all arcs and using additional assumptions already outlined in Sect. 4.1, we arrive at a cumulate recycling rate of  $1.5\text{--}6 \text{ km}^3/\text{year}$  ([4]), which is not insignificant when compared to the total recycling rate of  $\sim 20 \text{ km}^3/\text{year}$  for oceanic crust (basaltic + gabbroic sections [61]). However, it is not clear whether primary

garnet pyroxenite accumulation occurs in all arcs ([16, 23]). Island arcs and young, incipient continental arcs (e.g., the Cascades) appear not to have a strong garnet pyroxenite signature (compare Cascades data to Sierran data in Fig. 1), indicating that garnet pyroxenite was probably never on the liquidus. The lack of a garnet pyroxenite signature in young arcs and island arcs may stem from the fact that the preexisting lithosphere through which mantle-derived magmas pass is thinner in these regions [94]. Under these conditions, crystallization occurs at lower pressures where olivine is the dominant liquidus phase. If primary garnet pyroxenites only accumulate in continental arcs, then our calculated delamination recycling flux of garnet pyroxenites is a maximum bound. However, it is still possible for garnet pyroxenites to form in island arc environments by the subsolidus conversion of lower pressure cumulates (gabbros) to garnet pyroxenites. Once converted, these metamorphosed cumulates would be dense and could potentially delaminate as well. The only difference would be that garnet pyroxenites in continental arcs are primary cumulates, whereas those in island arcs might be metamorphosed low-pressure cumulates.

Given the possibility that the recycling flux of garnet pyroxenites via delamination might be significant, we end with the question of whether recycled arc-type garnet pyroxenites can be detected, that is, are these pyroxenites ever sampled in hotspots? The Sierran pyroxenites have the following primitive mantle-normalized ratios (denoted with subscript  $N$ ):  $(\text{Sm}/\text{Nd})_N \leq 1$ ,  $(\text{Lu}/\text{Hf})_N \leq 1$ ,  $(\text{U}/\text{Pb})_N$  only slightly greater than 1, and  $(\text{Rb}/\text{Sr})_N$  ranging between 1 and that for continental crust [18, 19]. In particular, these low  $\text{U}/\text{Pb}$  signatures may also be related to the fact that these garnet pyroxenites contain primary cumulate sulfide [42, 95]. These parent–daughter ratios will lead to time-integrated Nd, Hf, Pb and Sr isotopic compositions similar to the EM1 isotopic component seen in some hotspot magmas [96] and are hence reasonable candidates for certain recycled reservoirs believed to exist in the mantle [26]. Because of the likely lower melting temperatures of both the high MgO and low MgO Sierran garnet pyroxenites [97–100], such pyroxenite blobs could be preferentially melted and reacted with the mantle, forming the source reservoirs of some hotspot magmas [99, 101, 102]. Detecting these heterogeneities in the source regions of intraplate magmas will be aided by systematic study of major element [101, 103] and mildly incompatible element compositions (namely the first-row transition metals) [104–106].

## 5.2 Implications for Cretaceous to present volcanism in eastern Asia and western North America

There is growing evidence that arc-related pyroxenitic roots were generated beneath much of westernmost North

America during the Cretaceous and early Paleogene. Some of these mafic roots may have already foundered (as evidenced by seismic studies indicating hot asthenospheric mantle at depths where pyroxenites were once present based on the xenolith record [10, 11]), but some may still persist [12]. Is it possible that some of the volcanism associated with Basin and Range extension is related to preferential melting of these pyroxenite layers, as suggested long ago by Leeman and Harry [108]? Much work has been placed recently on estimating mantle potential temperatures of intraplate basalts [68, 83] in an attempt to constrain the thickness of the lithosphere and temperature of the asthenosphere. A fundamental step in estimating mantle temperatures from basalts is to back-correct for olivine fractionation until the liquid is in equilibrium with peridotitic mantle. However, if pyroxenites are interleaved with the peridotites, then smaller fractionation corrections are required and estimated temperatures would be lower. There is no easy way to resolve these differences.

In this context, it is useful to turn to the extensive studies of Cretaceous to Neogene basalts in eastern Asia, particularly in northeastern China. Much of eastern Asia was marked by continental arc volcanism in the Jurassic and Cretaceous, much like western North America at the same time [109–111]. Such volcanism must have left behind great thicknesses of mafic pyroxenites at the base of the crust or in the lithospheric mantle throughout much of eastern China. A curious feature of many basalts in eastern China is that they have high total Fe contents, which if one were to fractionate correct to nominal mantle peridotite would yield temperatures 100–200 °C higher than ambient mantle beneath ridges. However, a number of studies have noted that these basalts have unusual Fe/Mn and Zn/Fe systematics, which suggest a pyroxenite-bearing source rather than a pure peridotite source [106, 112–114]. It has been suggested that these pyroxenite-bearing mantle sources could be derived by re-melting of previously delaminated lower crustal pyroxenites (presumably arclogites?) [115, 116]. However, it is also possible that these pyroxenite lithologies may still be within the lithosphere, preferentially melting by extensional decompression if lithospheric thinning occurred by extension or thermal erosion, rather than by delamination (cf. [117–120]). In any case, the unusual transition metal systematics of Cenozoic basalts in eastern China (in the North and South China Blocks) suggest that the mantle melting regime beneath this area is more fertile. Recent work on peridotite xenoliths from the South China block seem to confirm the unusually high fertility of the mantle lithosphere [121]. No such detailed study of transition metal systematics has been conducted for western North American basalts, so it is unclear whether much of the volcanism in the area is driven by thermal variations in the asthenosphere or to the

presence of fertile heterogeneities associated with pyroxenite veins. A promising future direction of research will be to use basalts to interrogate both mantle fertility and temperature beneath eastern Asia and western North America.

## 6 Conclusions and future directions

In summary, we show that arc magmatism, particularly in continental arcs, generates thick and dense garnet pyroxenite cumulates, which eventually founder back into the mantle. The entire process, from initiating of magmatism to delamination, may take 10–30 My, implying that during the lifespan of a magmatic arc, numerous pulses of cumulate foundering may occur. What we have not yet considered are other processes that modulate the thickness of arc crust, such as erosion and tectonics (shortening and extension). Erosion and tectonics are also affected by magmatic inflation and vice versa, so the full feedbacks must also be considered. Another feedback that we did not consider was the effect of crustal thickening on suppressing the efficiency of decompression melting in the mantle wedge [55], which would influence crustal growth rates. Although our approach is simplistic and nature is complex, it is hoped that simple analytical solutions may provide a means of systematically evaluating each of these feedbacks in future studies.

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**Conflict of interest** The authors declare that they have no conflict of interest.

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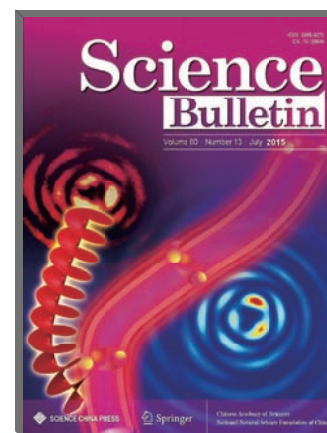
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**COVER** Accelerating optical beams have attracted a great deal of interest in the past several years, fueling the research in beam synthesis as well as other areas beyond optics. An intriguing question arises naturally: can we design accelerating beams that propagate along arbitrary trajectories and yet have controllable symmetric transverse profiles? In this issue, Chen's group provides a comprehensive review of recent work on various spatially-shaped accelerating beams along pre-designed trajectories, including self-accelerating, self-breathing, and self-propelling Bessel-like beams. They also discuss the potential application of these dynamical beams in optical trapping and manipulation. The cover figure illustrates an example of a specially-designed accelerating beam curving and propelling along a three-dimensional trajectory in free space (see the review by Juanying Zhao et al. on page 1157).



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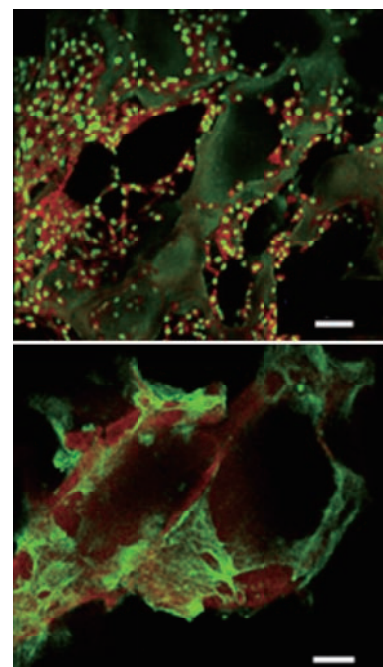
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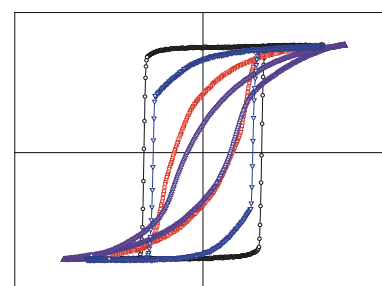
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